## Atomistic Models for Nanotube Device Electrostatics

## Slava V. Rotkin<sup>1</sup>, K.A. Bulashevich<sup>2</sup>, N.R. Aluru<sup>1</sup>

1 Beckman Institute for Advanced Science and Technology, UIUC, 405 N.Mathews Av., Urbana, IL 61801, USA E-mail: rotkin@uiuc.edu

2 St. Petersburg State Technical University, 21 Politekhnicheskaya, St. Petersburg, Russia

The 1D character of the electromagnetic eigenmodes of SWNT system results in a weak screening of the Coulomb interaction and the external field. Below we present a quantum mechanical calculation of the polarizability of the metallic [10,10] tube. This quantity is not defined solely by the intrinsic tube properties. In contrast, it depends also on the geometry of the nanotube, and/or the closest gates/contacts. Hence, the charge distribution has to be considered selfconsistently. The local perturbations of the electronic density will influence the entire system unlike in common semiconductor structures.

Figure 1 is a sketch of the depolarization of the tube potential (induced charge density) by the side electrode and the back-gate. The green continuos line is the statistical approximation (Boltzmann-Poisson eqs.) which coincides well with the Quantum Mechanical result (blue dotted line) except for the quantum beating oscillations at the tube end. The depolarization manifests as a significant nonuniformity of the charge along the tube length.

In one-dimensional systems, the potential,  $\phi^{act}$ , induced by a charge density,  $\rho^{ind}$ , is proportional to this charge density. Thus, the Poisson equation is effectively reduced to

$$\rho^{\mathrm{ind}}(z) = -e^2 \nu_M \phi^{\mathrm{act}}(z).$$

Here  $v_M$  stands for the nanotube density of states. We demonstrated that this quantity works as an atomistic capacitance of a SWNT:

$$C_A^{-1} = \frac{1}{e^2 \nu_M}$$

and

$$C_m^{-1} = 2\log\left(\frac{2h}{R}\right)$$

gives the geometric capacitance which is a function of distance to the back-gate and SWNT radius, in case of straight SWNT it is a logarithm. In equilibrium, we write a relation between the equilibrium charge density and external potential (gate voltage), which comprise both the atomistic and geometric capacitance:

$$\rho_{\infty} = -\frac{\varphi^{\text{xt}}}{C_m^{-1} + C_A^{-1}} \simeq$$

$$\simeq -\varphi^{\text{xt}} C_m \left(1 - \frac{C_m}{C_A}\right)$$

The last equation is still valid for the nanotube of an arbitrary shape, although, no simple expression for the geometric capacitance  $C_{\rm m}$  can be written for a nanotube when it is bent.

Authors acknowledge support through a CRI grant of UIUC (S.V.R., N.R.A.), DoE grant DE-FG02-01ER45932 (S.V.R.), RFBR grant 00-15-96812 (S.V.R., K.A.B.) and Beckman Fellowship from the Arnold and Mabel Beckman Foundation (S.V.R.).

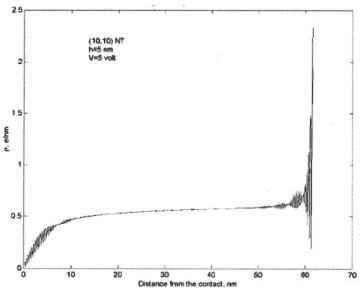


Figure 1. Selfconsistent charge density of a [10,10] armchair nanotube at 5 V voltage applied between side and back-gate contacts. The distance between the tube center and the back-gate is 5 nm, the tube radius and length are  $0.6\,\mathrm{nm}$  and  $60\,\mathrm{nm}$ .